

The MVACS Tunable Diode Laser Spectrometers

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Abstract. Two independent tunable diode laser spectrometers are resident aboard the Mars Polar Lander as part of the Mars Volatile and Climate Surveyor payload. One spectrometer is located on the meteorological mast for measurements of H₂O and CO₂ in the free atmosphere, and the other serves as the H₂O and CO₂ analyzer for the thermal and evolved gas analyzer. Water vapor is measured using a diode laser operating at 1.37 μm , while CO₂ is measured using a second laser operating near 2.05 μm . The 2.05 μm laser also has isotopic analysis capability. In addition to the major CO₂ isotopomer (¹²C¹⁶O¹⁶O), analysis of ¹³C¹⁶O¹⁶O and ¹²C¹⁸O¹⁶O in the atmosphere and in the thermal and evolved gas analyzer are possible under certain conditions. The spectrometers were designed and built at the Jet Propulsion Laboratory and have their heritage in a series of tunable diode laser spectrometers developed for Earth atmospheric studies using high-altitude aircraft and balloon platforms. The 1.37 μm TDL on the meteorological mast will provide the first in situ measurements of water vapor in the martian boundary layer, with a sensitivity nearly two orders of magnitude greater than the water vapor abundances inferred from the remote sensing observations by the Viking Orbitors.

1. Introduction

Tunable diode lasers (TDLs) have been used in high resolution spectroscopic studies of gases since the 1970s. Their extensive wavelength coverage (0.5 - 30 μm), rapid and continuous tunability, and narrow spectral linewidth make them ideal as light sources in gas sensing spectrometers. Lead-salt TDLs cover the wavelength region from 3.2 μm to beyond 30 μm , but require cooling to temperatures typically below 120 K which prohibits their use on an extended planetary mission such as the Mars Polar Lander. TDLs that operate continuous wave at wavelengths in the 0.6 to 2.1 μm region at much higher temperatures (up to 50 C) are available as a result of extensive development by the optical communication industry. That industry requires single-frequency lasers operating at the near-infrared wavelengths of 1.3 and 1.55 μm where optical fiber attenuation coefficients are at a minimum. Early work at JPL [Forouhar *et al.*, 1992] extended the wavelength range of single-frequency, high-temperature (300 K) devices to near 2.1 μm , which enables a wide range of gases to be monitored using overtone or combination bands present in the near-infrared region.

H₂O and CO₂ both have relatively strong absorption bands in the 1-2 μm region. The $\nu_1 + \nu_3$ combination band of H₂O is centered near 1.38 μm , while three bands of CO₂ and its isotopomers occur in the 2 μm spectral region. Isotopic bands of these molecules will be shifted in wavelength relative to the primary isotope. A careful analysis of the spectra of H₂O and CO₂ was undertaken to select the optimum measurement regions for the MVACS TDL spectrometers. One TDL spectrometer is located on the meteorological (MET) mast, and the other is part of the thermal and evolved gas analyzer (TEGA). For MET the goal was to select a CO₂ spectral region containing as many isotopic lines as possible while simultaneously avoiding the very strongest main isotope lines which would be too optically deep in the Martian 95% CO₂ atmosphere. For TEGA, the emphasis was placed on obtaining a low detection threshold for the main CO₂ isotopomer while maintaining the ability to monitor the carbon-13 and oxygen-18 isotopomers. Although HDO measurements are technically possible if the water vapor abundance is sufficiently large and long-term spectrum averages are performed, detection of HDO is marginal due to the extremely weak absorption levels expected in the 1.37 μm region.

Other papers in this special issue address the scientific rationale for the measurements of water vapor and carbon dioxide, and isotopic composition on Mars, and provide detailed descriptions of the MET and TEGA instruments (Boynton *et al.*, this issue; Crisp *et al.*, this issue). In this paper, the technical details of the two MVACS TDL spectrometers are presented, along with brief descriptions of the calibration and data processing methods which are described in detail elsewhere [May and Webster, 1993; May, 1998].

2. Heritage

The TDL sensors developed for MVACS have their heritage in a series of TDL spectrometers developed at JPL beginning in the early 1980s. At that time, diode lasers required cooling to temperatures in the 20-60 K range using liquid helium dewars, and field instrumentation was necessarily physically large and expensive. In 1985 the first mid-IR (3 - 10 μm) TDLs became available that operated at 77 K, and slightly above, enabling a significant reduction in instrument size as liquid nitrogen could replace liquid helium as the cryogen in the laser dewars. A balloon-borne, multi-laser TDL spectrometer (BLISS) was flown successfully twelve times between 1982 and 1992 [Webster and May, 1987; Webster *et al.*, 1990] and made

measurements of NO, NO₂, HNO₃, O₃, CH₄, N₂O, HCl, H₂O, HDO, and CO₂ in the Earth's atmosphere. Both liquid helium and liquid nitrogen dewars were utilized on BLISS, and the instrument package had a mass of nearly 1000 kg.

The success of the BLISS instrument led to the development of a ALIAS, a much smaller, four-channel TDL spectrometer for the NASA ER-2 aircraft [Webster *et al.*, 1994] which used only lasers that could operate at temperatures above 77 K. This instrument, and a two-channel version developed for a balloon platform (ALIAS-II, Scott *et al.* 1999) are both currently operational and participating in NASA field campaigns. Each houses the TDLs in a liquid nitrogen dewar with a typical operational time between cryogen refilling of >12 hours (30 hours non-operational). ALIAS-II has a mass of 28 kg which, although small relative to prior configurations, is still a research-grade field spectrometer requiring daily cryogen replenishment. In 1990 the PIRLS instrument was proposed for the NASA Cassini mission by Webster *et al.*, 1990. This instrument design utilized a novel Joule-Thomson cooler developed at JPL to cool lead-salt TDLs to near their nominal 80 K operating temperature, and was also equipped with a particle size spectrometer. It was the first comprehensive effort at designing and breadboarding a space-qualifiable TDL spectrometer at JPL. Total instrument mass for PIRLS was approximately 5 kg.

At present, diode lasers operating continuous wave at wavelengths longer than 2.1 μm must be cryogenically cooled, and many gases of interest can only be monitored at low levels (parts-per-billion, and below) using mid-IR lasers which can access the fundamental vibration bands. However, water vapor has a relatively strong absorption band centered at 1.38 μm which is accessible using TDLs that do not require cryogenic cooling. CO₂ and several of its isotopes can be measured at 2.05 μm using similar lasers. In 1992 near-IR (1 - 2 μm) diode lasers were developed at JPL and optimized during the next three years for use in a new class of TDL spectrometer. These devices produced single-frequency output as a result of the distributed feedback (DFB) structure [Forouhar *et al.*, 1993], and provided several mW of continuously-tunable output. In 1996 the first TDL-based JPL aircraft instrument for measurements of water vapor from the NASA high-altitude ER-2 research aircraft was flown [May, 1998], and in 1998 a version for the NASA DC-8 was deployed to make measurements in thunderstorms and hurricanes as part of the third Convection and Moisture Experiment (CAMEX-3). These near-IR aircraft instruments formed the basis for the design of the MVACS spectrometers, which incorporate many of the same basic design features.

Substantial reductions in mass, volume, and power consumption were necessary for integration into the MVACS payload. Initially, only measurements of water vapor in the free atmosphere were to be undertaken. As the instrument design phase progressed, a second laser was added with the goal of analyzing the abundances of isotopic CO₂ in the atmosphere, and eventually a dual laser system was also selected as the evolved gas analyzer for the TEGA. Volume and power constraints led to significant design compromises that resulted in reduced sensitivity for the MVACS TDL spectrometers in comparison to those developed at JPL for aircraft and balloon platforms. However, the primary science goals of H₂O and CO₂ detection in the Mars atmosphere, and in evolved gas samples, can be met with the flight hardware. In the following sections the instrument design details are given along with specifications and expected performance for the MET and TEGA TDL spectrometers.

3. Instrument Descriptions

The MET and TEGA TDL spectrometers share common designs for several subsystems. Among these are the control and signal processing electronics (jumped for slightly different operation on the two systems), the laser and detector optical mounting assemblies, and the use of multipass (Herriott) absorption cells used to increase the optical absorption pathlength. Differences between the MET and TEGA systems include the Herriott cell mirror sizes and focal lengths, minor thermal control details, and the exact laser wavelengths used to monitor the chosen H₂O and CO₂ rovibrational lines. Table 1 summarizes the spectrometer specifications.

Table 1. TDL Spectrometer Specifications

Parameter	MET	TEGA
Mass (excluding electronics, kg)	0.5	0.44
Power Consumption (maximum, W)	4.3	4.5
Optical Pathlength (cm)	1055.6	100.8
Mirror separation (cm)	27.78	5.60
Mirror diameter (cm)	3.40	2.54
Mirror center thickness (cm)	0.5	0.4
Mirror substrate material	Zerodur	Zerodur
Mirror optical coating	Cr-Au, Al ₂ O ₃ overcoat	Cr-Au, Al ₂ O ₃ overcoat
Number of beam passes	38	18
Inner coupling hole radius (cm)	0.90	0.47
Outer coupling hole radius (cm)	1.30	.87
Detector (active area diameter)	InGaAs (1 mm)	InGaAs (1 mm)
H ₂ O laser set point (C)	22.8	2.9
CO ₂ laser set point (C)	17.9	-3.5
Beam injection angle (inner hole, deg)	1.703	3.974
Beam injection angle (outer hole, deg)	2.460	7.356
Nominal beam diameter (all lasers, cm)	0.25	0.25
H ₂ O spectral region (cm ⁻¹)	7297.9 - 7300.4	7306.0 - 7307.5
CO ₂ spectral region (cm ⁻¹)	4885.0 - 4886.9	4876.7 - 4878.6
Minimum detectable absorption level	4 x 10 ⁻⁴	8 x 10 ⁻⁴
H ₂ O minimum detectable abundance	5.6 x 10 ¹¹ cm ⁻³	1.0 x 10 ¹³ cm ⁻³
CO ₂ minimum detectable abundance	4.2 x 10 ¹⁴ cm ⁻³	7.2 x 10 ¹⁴ cm ⁻³

Lasers

Distributed feedback (DFB) lasers operating at 1.37 μm and 2.05 μm (*Forouhar et al*, 1992; *Forouhar et al*, 1993) were fabricated and flight-qualified at JPL for the MVACS TDL spectrometers. These devices operate on a single longitudinal mode with a side mode suppression of >30 db, so spectral interference from minor modes is negligible. The laser chips are mounted in 5.6 mm cylindrical "TO" packages and hermetically sealed in the packages with welded endcaps containing anti-reflection (AR) coated windows. Output power is in the 5-8 mW

range for the 1.37 μm lasers, and 2-3 mW for the 2.05 μm lasers. Continuous tuning ranges exceeded 2 cm^{-1} for all devices.

Electronics

TDL Control and signal processing electronics are contained on a single, six-layer circuit board measuring 8.4 x 19.1 cm, which is mounted in the payload electronics box (PEB). The circuit consists of two functional sections for each laser/detector channel that are associated with control of the laser current, and detector signal processing. A block diagram of the TDL circuitry is shown in Fig. 1. Spectra are recorded by holding the laser temperature fixed, and changing the wavelength by varying the injection current. The laser current is swept repetitively over a range of a few tens of mA which corresponds to a spectral region of about 2 cm^{-1} . A novel design for controlling the laser scan waveform was developed for MVACS and is described in detail by *Woodward and May, 1999*. A programmable ROM holds two laser scan waveforms that are continuously "played out" to the laser drive circuitry via logic that cycles through the ROM addresses and delivers a programmed laser current value to the appropriate laser. A small-amplitude (~ 0.5 mA peak) sine wave is superimposed onto the base laser scan ramp to enable second harmonic (2f) detection. Both the ramp and sine wave amplitudes are programmable, as is the ramp beginning and end points which define the spectral scan region. For the MVACS TDL spectrometers the laser scan period is fixed at 2.79s, limited by the rate at which data could be transmitted across the 9600 baud serial link to the lander computer. The laser modulation frequency is 2.93 KHz (2f detection frequency 5.86 KHz).

Detector signals are provided as inputs and are processed using standard second harmonic detection techniques [*Webster et al., 1988; May and Webster, 1993; May, 1998*]. Briefly, both the dc and ac components of the detector signal are extracted and downlinked for processing. The dc component represents the laser power and is required for normalization of the ac signal. The ac component is the second harmonic spectrum obtained by demodulation of the detector signal at 5.86 KHz (twice the laser modulation frequency). Fig. 2 shows examples of dc and 2f spectra recorded by the TEGA TDL for a pure CO_2 sample. The laser power increases with injection current which produces the sloping baseline in the dc spectrum. Since the 2f signal amplitudes are proportional to the received laser power at the detector, the dc spectrum must be used for normalization before the 2f signals can be used for quantitative determination of gas concentration. No spectrum processing is performed by the lander computer.

Optics

Both the MET and TEGA spectrometers utilize multipass "Herriott" cells [*Herriott et al., 1964; Altmann et al., 1981*]. These configurations produce a 1.0 m optical absorption path for the TEGA spectrometer, and a 10.6 m absorption path for the MET spectrometer. For each Herriott cell the laser beams are coupled into the system via holes in one mirror (the "coupling" mirror) that lie on different radii so that two nonoverlapping spot patterns are formed on the mirror pairs. The beams exit the system through the input coupling holes and impinge upon detectors which are mounted on the same thermal plate as the lasers. For each system the two laser beams make nearly circular spot patterns on the Herriott cell mirrors with radii defined by the coupling hole positions. Both mirrors in a Herriott cell have the same radius of curvature. Zerodur was chosen as the mirror substrate material for both the MET and TEGA systems, with

Cr-Au as the reflective coating. An Al_2O_3 overcoat protects the Au surfaces. This coating was found to perform extremely well with the external optics flown on the NASA ER-2, WB57F, and DC-8 aircraft experiments, which were subjected to >200 m/s airflow continuously during flights. No significant reduction in mirror reflectivity was observed after more than 500 flight hours. Thus, for the MET system, we do not expect mirror damage from dust or aerosol impact for typical wind conditions on Mars.

Fig. 3 shows a single laser/detector assembly (for the TEGA $2.05\ \mu\text{m}$ laser) which illustrates the basic design employed for laser and detector mounting for all four MVACS laser/detector heads. A single aspheric lens is used to collect and collimate the laser output beam. This arrangement produces a collimated ellipse with a major/minor axis ratio of approximately three. The lens is positioned to achieve a beam focus that is slightly beyond the midpoint between the Herriott cell mirrors. This results in a beam spot size at the detector that is optimum for the 1 mm detector area without the need for a second collecting/focusing lens at the detector.

Two angled aluminum (Al) mounts hold the laser packages at the appropriate Herriott cell injection angles. The laser collimating lens is a 4.5 mm diameter asphere with 4.5 mm focal length, and is mounted to the same Al piece that holds the laser package. Each detector is a 1 mm diameter InGaAs photodiode operated with zero bias voltage. Extended-range InGaAs detectors are used for the $2.05\ \mu\text{m}$ lasers. The laser/detector separation is a function of distance from the Herriott cell coupling mirror as determined by the composite angle between the entrance and exit beams. A 6.6×6.6 mm, 1.72 W, thermoelectric cooler (TEC) is sandwiched between the Al laser mount and a second Al bracket which serves both as the “sink” for the TEC, and the supporting mount for the laser/detector assembly. Each TEC operates as a cooler or a heater, as required, but will function primarily as a heater at Mars because the laser set points (near 0 C for TEGA and near 20 C for MET) are above the expected ambients during most of the mission. Both the laser and detector temperatures are held constant within ± 0.01 C using hybrid controller circuits provided by Hytek Microsystems. These controllers were configured for true proportional-integral-derivative (PID) operation to minimize overshoot on power up. To further improve laser temperature setpoint stability against excursions of Mars ambient, MET TEC controller circuitry includes provision for electronic compensation of the effects of finite thermal conductivity of the Al mount block. This novel arrangement dynamically adjusts controller setpoint as a function of TEC power input to anticipate and partially cancel temperature control errors resulting from thermal gradients within the laser mount assembly.

MET

The MET spectrometer consists of an *open-path* Herriott cell mounted on the MET mast, and two diode lasers for measurements of atmospheric H_2O and isotopic CO_2 . Fig. 4 shows the layout on the MET mast. The optical absorption pathlength for both MET TDL channels is 10.6 m. The multipass mirrors are mounted in fixed locations on the mast and exposed to the free atmosphere. Temperature and pressure data, required for TDL spectrum processing, are supplied by independent sensors whose details are described by *Crisp et al* (this issue). Due to sequencing and data acquisition constraints, spectra of H_2O and CO_2 are not acquired simultaneously but are recorded alternately with only one of the two lasers operating at any given time.

In Fig. 5 example spectra are shown from single laboratory scans of each MET laser (2.79 s acquisition time). These spectra were recorded with the lasers mounted in custom laboratory test fixtures which allowed the laser operating temperature and current to be adjusted easily. In addition, the necessary alignment adjustments were available to minimize optical fringing in the spectra. In the flight systems there was insufficient movement in the collimating lens position in the plane perpendicular to the optic axis of the lens to achieve optimum alignment. Therefore, optical fringing levels in the spectra are significantly higher than they would be for a properly aligned system and are the limiting noise factor by a full order of magnitude. Base line center absorption levels are in the 5×10^{-4} range, which can be translated into equivalent minimum detectable number densities using Beer's law. The results are listed in Table 1. The effects of optical fringing can be reduced somewhat by fitting and analysis procedures, as well as spectrum averaging. Fig. 6 shows an example of fringe reduction for a single MET CO₂ spectrum. Such techniques will be required to extract maximum information from the MVACS TDL spectrometers due to the excessive amplitude of the optical fringes in the spectra. For comparison, the Viking MAWD instruments [Farmer *et al.*, 1979] measured minimum water vapor abundances of approximately one precipitable micron, which if uniformly mixed with pressure, would yield surface water vapor concentrations on the order of 3.0×10^{12} .

TEGA

The TEGA spectrometer is a *closed-path* system where evolved gases are brought into the Herriott absorption cell via heated transfer lines and a stream of N₂ carrier gas. Details of the TEGA instrument are given in Boynton *et al.* (this issue). The multipass mirrors are mounted within an insulated Al tube. Fig. 7 shows the layout of the TEGA TDL spectrometer and indicates the positioning of the laser/detector assemblies with respect to the Herriott cell mirrors. Note that a portion of the optical path includes the area within the optical head between the laser and the Herriott cell coupling mirror, and between this mirror and the detector. Quartz windows attached to the back of the coupling mirror isolate the internal Herriott cell volume (and optical pathlength) from the optical head, but the optical head is not purged and will contain ambient mars atmosphere. Absorption within this portion of the optical path will be taken into account by recording spectra while pure N₂ purge gas is flowing through the Herriott cell, and subtracting the resulting spectrum from the spectra recorded during the subsequent heated sample run.

Example spectra recorded with the flight unit are shown in Fig. 8. Unlike the MET system, both lasers are operated during sample runs for the TEGA so that spectra of H₂O and CO₂ can be acquired simultaneously. As with the MET spectrometer, optical fringes limit the sensitivity due to absence of necessary optical adjustments for the TDL collimating lens. Conventionally, IR spectra are displayed in wavenumber (cm⁻¹) units, with wavenumber increasing from left to right as shown. Because the near-IR diode lasers tune such that wavenumber decreases (wavelength increases) with increasing current, the spectra shown in Fig. 8 (and Fig. 5) are reversed in comparison to the raw spectral data that are downlinked (eg. compare the peak positions in Figs. 2 and 8).

3. Calibration and Data Analysis

Second harmonic detection is implemented for the MVACS spectrometers. Details of this detection technique, and data interpretation methods, can be found in Webster *et al.*, 1988,

May and Webster, 1993, and *May*, 1998. Briefly, a constant-period sawtooth current ramp is applied to the TDL to scan the output wavelength over the desired spectral interval (typically 1-3 cm^{-1}). A small-amplitude sinusoidal waveform at frequency f is added to the ramp, and the detector signal is demodulated at $2f$ to produce the second harmonic spectrum. The amplitude of the sinusoidal modulation is adjusted to produce the optimum $2f$ spectrum based on the conditions of the measurement. Since the molecular line shape varies with pressure and temperature, some care is required in choosing the optimum modulation amplitude if its value is not continually updated by the software based upon current pressure and temperature measurements. For MVACS the modulation amplitude is held fixed for all lasers at values optimum for 7 mbar pressure (MET) and 15 mbar (TEGA) because the modulation amplitude is programmed into the waveform ROM and cannot be changed during operations at Mars.

Molecular number densities are extracted from the observed second harmonic signal amplitudes using the Beer-Lambert law, and appropriate analysis of the second harmonic lineshapes. Full details are contained in *May*, 1998 and are therefore not repeated here. Calibration of the $2f$ signal chain is accomplished by comparing the direct absorption spectrum to the corresponding $2f$ spectrum and formulating a "response" number, R , for each signal chain. Details are described in *May and Webster*, 1993. Fig. 9 shows a modulated direct transmission spectrum (dc spectrum) for a pure CO_2 sample. A modulated direct transmission spectrum will accompany each MVACS $2f$ spectrum, and when the concentration of the target gas is sufficient to observe the spectral line in the direct transmission spectrum it will be possible to confirm the R number values from the actual Mars spectra. The direct transmission spectral line is split from the wavelength modulation which places the laser wavelength off of the line at the extremes of the cosine modulation function. The dashed line in the modulated direct transmission spectrum is a synthetic best-fit spectrum generated for the modulation amplitude used (0.035 cm^{-1} peak), and for the CO_2 pressure (5.73 mbar for this example) required to match the observed modulated direct transmission level. For each target spectral line, the peak-to-peak $2f$ signal amplitude (in counts from the A/D converter) is divided by the observed laser power at line center for that line (also in A/D counts, from the dc spectrum) to obtain the normalized $2f$ signal level. This process defines the R number for that spectral line which is essentially a measure of the $2f$ signal chain gain. Note that an accurate pressure measurement (for a pure gas, or a precisely known abundance of the target gas for an impure sample), and knowledge of the laser modulation amplitude, are required for determination of R for each channel. For Mars spectrum analysis, a data processing matrix describing the variation in $2f$ signal amplitude with total pressure and temperature is consulted to make the necessary corrections for the change in $2f$ lineshape with pressure and temperature (see *May*, 1998 for an example of such a matrix), and to extract the absolute $2f$ signal amplitude for a given number density of target gas.

Line strengths for the target H_2O lines have been measured by *Toth* (1994) using high-resolution Fourier transform spectroscopy. Extensive line strength measurements for a strong H_2O line in this same spectral region by *May*, 1998 are in excellent agreement (<2%) with *Toth's* value and support his absolute line intensities in this region. Net uncertainty in the H_2O line strengths is <4% for the TEGA and MET target lines. At typical Mars ambient pressures (<10 mbar) the pressure broadened component of the observed line widths for H_2O will be approximately 10% of the Doppler component. Therefore, uncertainties in the CO_2 broadening coefficient for H_2O will not have a significant impact on derived water vapor volume mixing ratios. For CO_2 , the line position and intensity data are taken from the 1996 HITRAN database

(Rothman *et al.*, 1998) for all isotopes. References there describe how the HITRAN data were assimilated for CO₂ in the 2 μ m region, and the expected uncertainties. Line strengths for the CO₂ spectral lines used for both the MET and TEGA spectrometers have estimated uncertainties below 5%. As with H₂O, pressure-broadening for CO₂ (self-broadening in this case) is relatively small at the low pressures expected at Mars. However, it represents a larger fraction of the observed linewidth for CO₂ due to the higher molecular weight, and the longer wavelength, which combine to reduce the Doppler component of the linewidth. The Doppler width for CO₂ at 273 K and 4880 cm⁻¹ is 0.0044 cm⁻¹ compared to a self-broadened component of approximately 0.001 cm⁻¹ at 10 mbar and 273 K.

Minimum detectable number densities for H₂O and CO₂ are listed in Table 1 for both the MET and TEGA TDL spectrometers. These values are based on measurements from single spectra with no spectrum averaging or fringe reduction efforts. Note that the spectral regions are different for the two spectrometers so the relationships between minimum detectable absorption level (in terms of fractional absorption) and minimum detectable gas number density are different. The base detectable line center absorption level in the TEGA spectra is 8×10^{-4} , while for the MET spectra the minimum detectable absorption level is 4×10^{-4} . These values are an order of magnitude larger than is relatively easily attainable in a properly aligned second harmonic TDL spectrometer, and result from the inability to position the collimating lens in the correct position for these systems. Since the fringing levels vary slightly with the measurement conditions (temperature primarily), the values listed are typical of those observed using the flight spectrometers and are based on peak-to-peak fringe levels in the spectra.

The ¹³C¹⁶O₂ and ¹²C¹⁸C¹⁶O isotopomers of CO₂ are present in both MET and TEGA spectral regions as shown in Fig. 5 and Fig. 8. Based on the minimum detectable absorption levels measured from single spectra, the precision to which the ¹²C/¹³C and ¹⁶O/¹⁸O isotopic ratios can be estimated if an assumption is made for the total number density of CO₂. For the MET spectrometer, assuming a 95% CO₂ atmosphere, 7 mbar total pressure, and 240 K temperature the precision in the isotopic ratios measured from single spectra is approximately 1-2% for both ¹²C/¹³C and ¹⁶O/¹⁸O since the ¹³CO₂ and ¹²C¹⁸O¹⁶O spectral lines have approximately equal intensity. Averaging small groups of spectra, and applying fringe reduction techniques, should improve the precision by at least a factor of two, and in some cases as much as a factor of four. Longer term spectrum averages (over many MET sessions) should lead to significant improvement in the precision of the isotopic ratios since the optical fringes in the spectra average out well over longer time periods. For TEGA, the ¹⁶O/¹⁸O ratio will probably not be measureable to better than 30% even for a pure CO₂ sample. The ¹²C/¹³C ratio should be measureable to approximately 10% for a pure CO₂ sample with no spectrum averaging, and to approximately 3-4% with extensive spectrum processing. These values can be linearly scaled for lower CO₂ abundances more typical of those expected during an actual TEGA run. Longer term spectrum averages are not feasible for TEGA due to the nature of the experiment (ie. the time dependence of the evolved gases). Again, excessive optical fringe levels in the spectra limit the precision to which isotopic ratios can be determined.

There is a very weak HDO line within the MET spectral region at 7298.1 cm⁻¹ that might be observable for sufficiently high levels of water vapor, optimum spectrum processing, and longer term spectrum averaging. The calculated line center absorption for this HDO line, using the HDO line parameters of *Toth*, 1994, is 4×10^{-5} for a total water vapor number density of 1×10^{14}

molecules/cm³ at 7 mbar total pressure and 270 K (this number density corresponds to approximately 500 ppmv) and an assumed D/H ratio that is six times greater than on Earth. Based on the single-spectrum minimum detectable line center absorption value of 4×10^{-4} this line would not be observable. However, with multiple spectrum averages over many MET sessions this line may be detectable and would provide a direct measurement of the atmospheric D/H ratio.

The MVACS TDL instruments are the first diode laser spectrometers to be designed, qualified, and flown on a planetary mission. Although significant design compromises were necessary to accommodate these instruments within the MVACS payload, their inherent small size, low power consumption, lack of moving parts, and high detection sensitivity for a wide range of gases makes TDL spectrometers viable and valuable tools for future space-borne and planetary instrument payloads. Improvements to the electronics and optical designs, which have already been implemented in prototypes currently being built at JPL for the International Space Station, along with new semiconductor laser sources such as quantum cascade lasers which can operate at wavelengths as long as 17 μm at room temperatures, should greatly improve instrument performance and utility for space-borne applications of semiconductor laser spectrometers.

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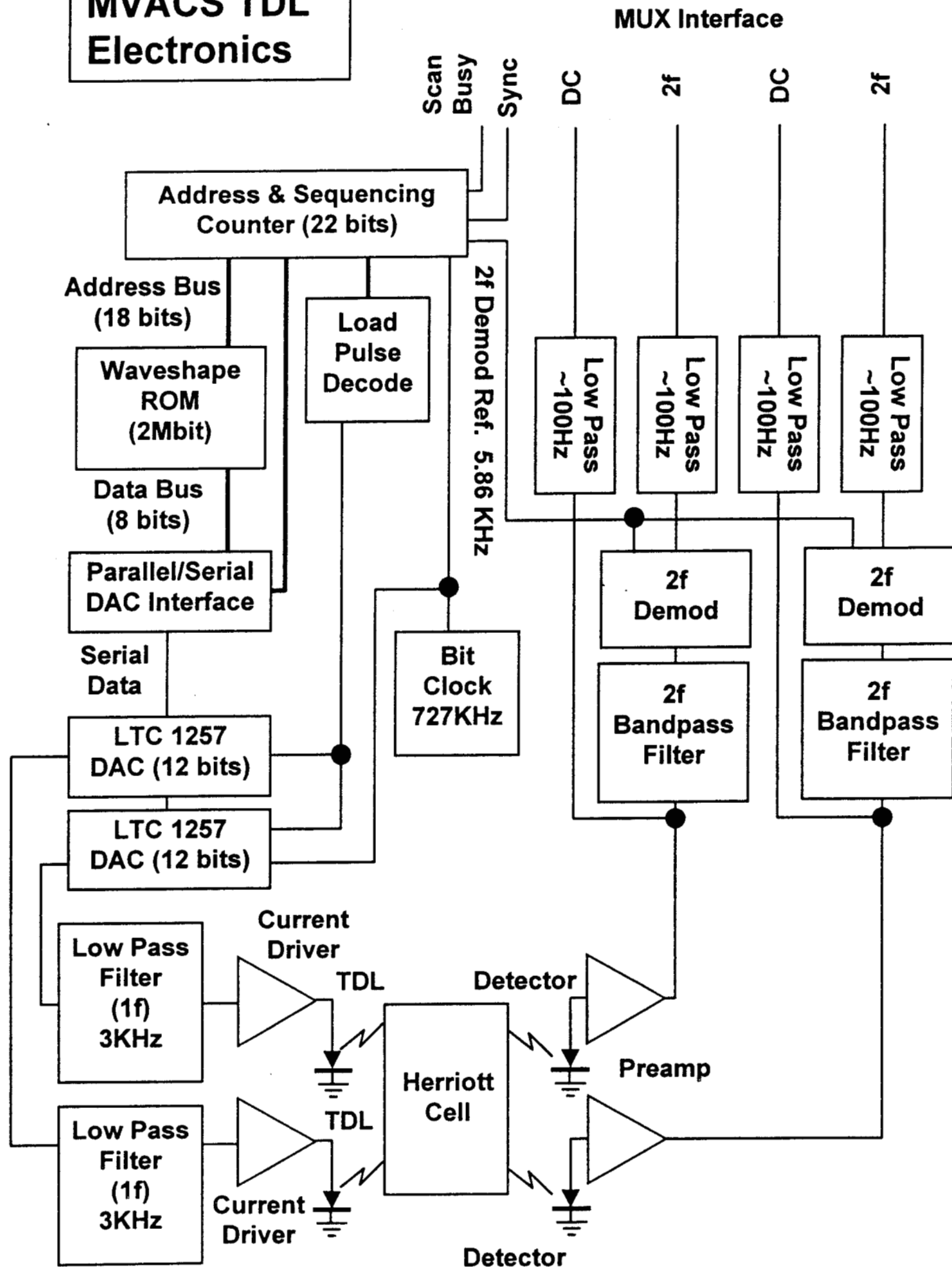
Figure Captions

1. Block diagram of the MVACS TDL control and signal processing electronics. Laser current scan waveforms are stored in the waveshape ROM and continuously played out to the laser current driver circuitry to sweep the lasers over the designated wavenumber range. The dc component of the detector signal is amplified and low-pass-filtered and supplied as an analog output signal. The ac component at 5.86 KHz is the second harmonic (2f) spectrum and is extracted using a bandpass filter, a synchronous demodulator, and a low-pass filter.
2. Example of raw TDL spectra from the TEGA spectrometer for a pure CO₂ gas sample at approximately 9 mbar pressure and 300 K temperature. The laser power (shown in A/D converter counts) increases with laser injection current, and the modulated CO₂ lines are visible in the upper spectrum. The lower spectrum is the 2f spectrum which must be normalized to the laser power spectrum for quantitative analysis. For this process a polynomial baseline is fit to the laser power spectrum and the value at the 2f line center peaks used for normalization of the 2f signal amplitudes.
3. Photograph of a laser/detector package for the TEGA 2.05 μm laser. The laser is mounted behind the collimating lens, which is visible as the upper circular optic. The detector is mounted directly below the laser and is not equipped with a collimating lens. Appropriate injection angles for the Herriott cells were machined into the mounting pieces, eliminating the need for external beam steering mirrors.
4. Photo of the MET TDL spectrometer as mounted on the MET mast. The optical heads are contained in an Al enclosure shown towards the rear of the photo with the two radiator disks on top. The face of the Herriott cell coupling mirror can be seen here, with the opposing Herriott cell mirror mounted in an Al mount located approximately 28 cm away (Table 1).
5. Spectral regions chosen for monitoring H₂O and CO₂ for the MET TDL spectrometer. The numeric labels in the lower CO₂ spectrum identify the isotopomers (626 = ¹²C¹⁶O¹⁶O, 636 = ¹³C¹⁶O¹⁶O, 628 = ¹²C¹⁶O¹⁸O). Signal amplitudes are in counts from the A/D converter.
6. Example of optical fringe reduction via an interactive software analysis method. The optical fringes arise from standing waves set up within the laser/detector path, and could not be further minimized for the MVACS TDL systems because of insufficient lateral adjustment range for the laser collimating lens. Signal amplitudes are in A/D counts.
7. Layout of the TEGA TDLs and Herriott cell. The laser/detector assemblies are located in a chamber that is connected to the insulated (aerogel) Herriott cell, but thermally isolated from the Herriott cell itself via a G-10 insulating washer. Gas enters the Herriott cell through a 1/16" O.D. stainless steel tube, and exits at the opposite end of the cell where a pressure sensor is located. The Herriott cell temperature is stabilized at 308 K to avoid condensation of H₂O.

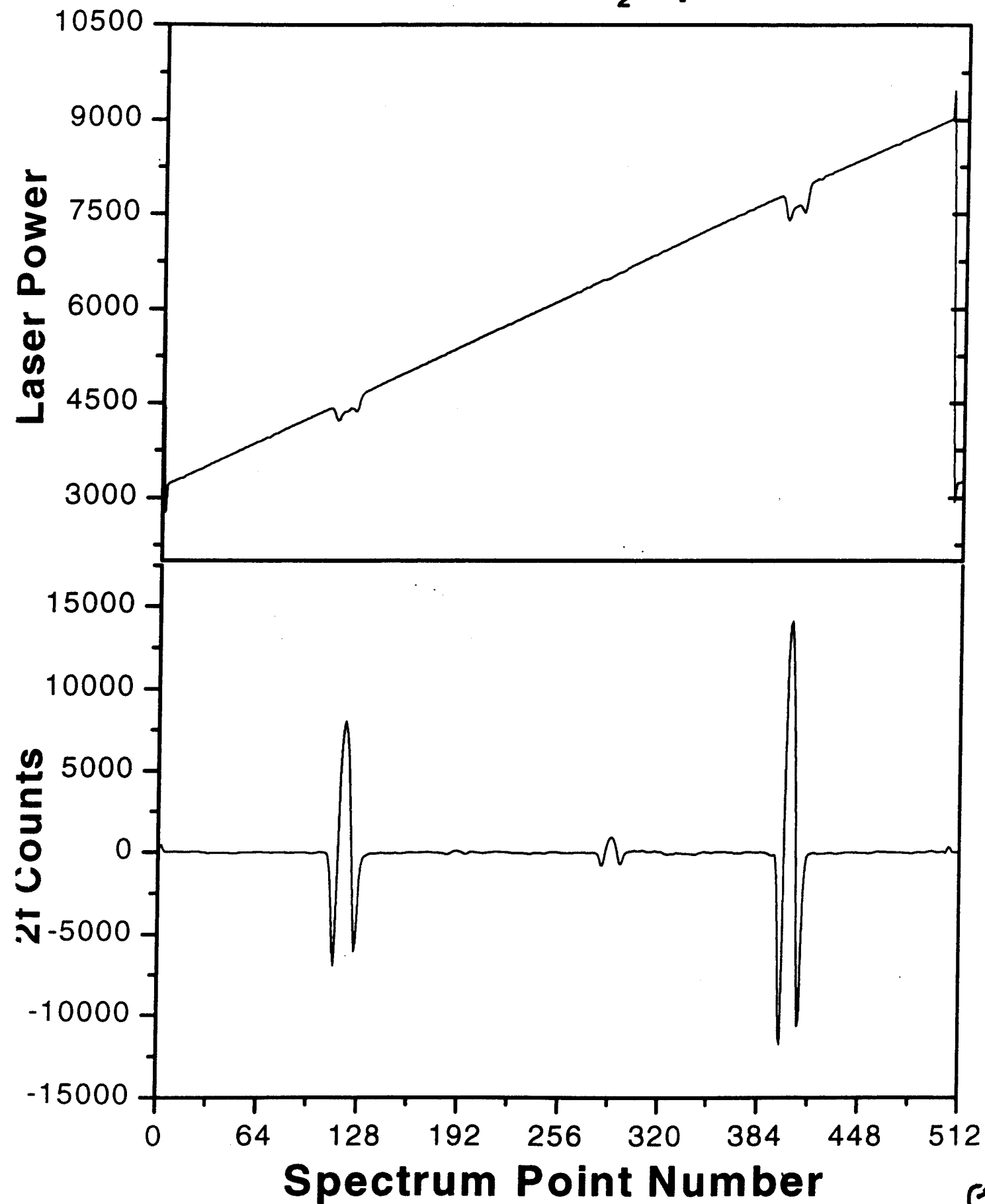
8. Spectral regions chosen for monitoring H_2O and CO_2 for the TEGA TDL spectrometer. The numeric labels next to the CO_2 2f lines identify the CO_2 isotopomers as in Fig. 5. A weak 628 line is also present in the spectrum near 4878.0 cm^{-1} , but this line will only be observable for very high levels of CO_2 in the TEGA cell. Signal amplitudes are in A/D counts.

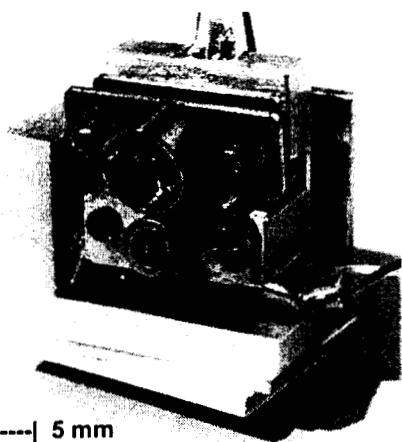
9. Modulated direct transmission spectrum (upper plot) and the corresponding 2f spectrum (lower plot) used to derive the 2f signal chain response number, R. Analysis of these spectra define the quantitative relationship between the observed 2f signal amplitudes and the gas number density for fixed electronic gain factors and laser modulation amplitude.

MVACS TDL Electronics



Raw TEGA CO₂ Spectra

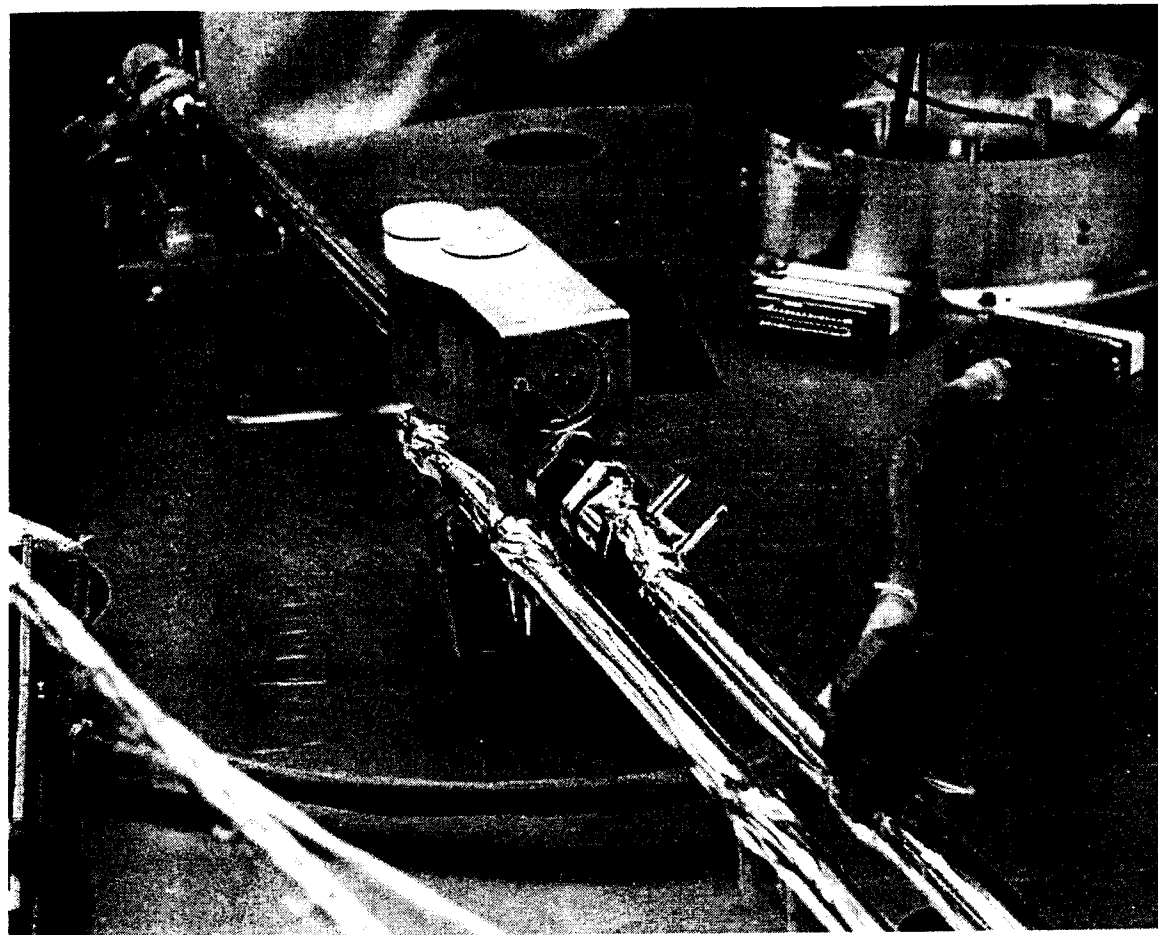




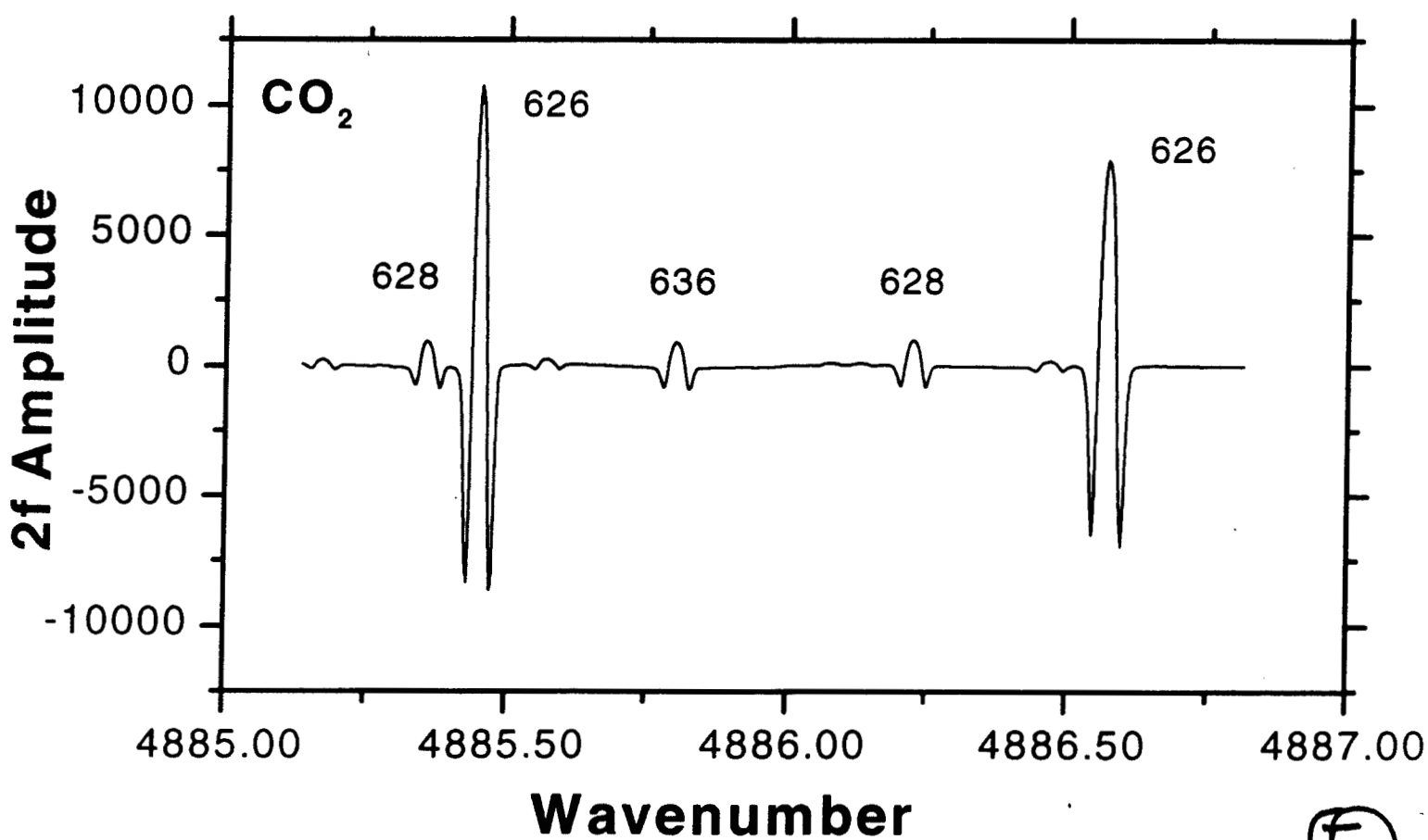
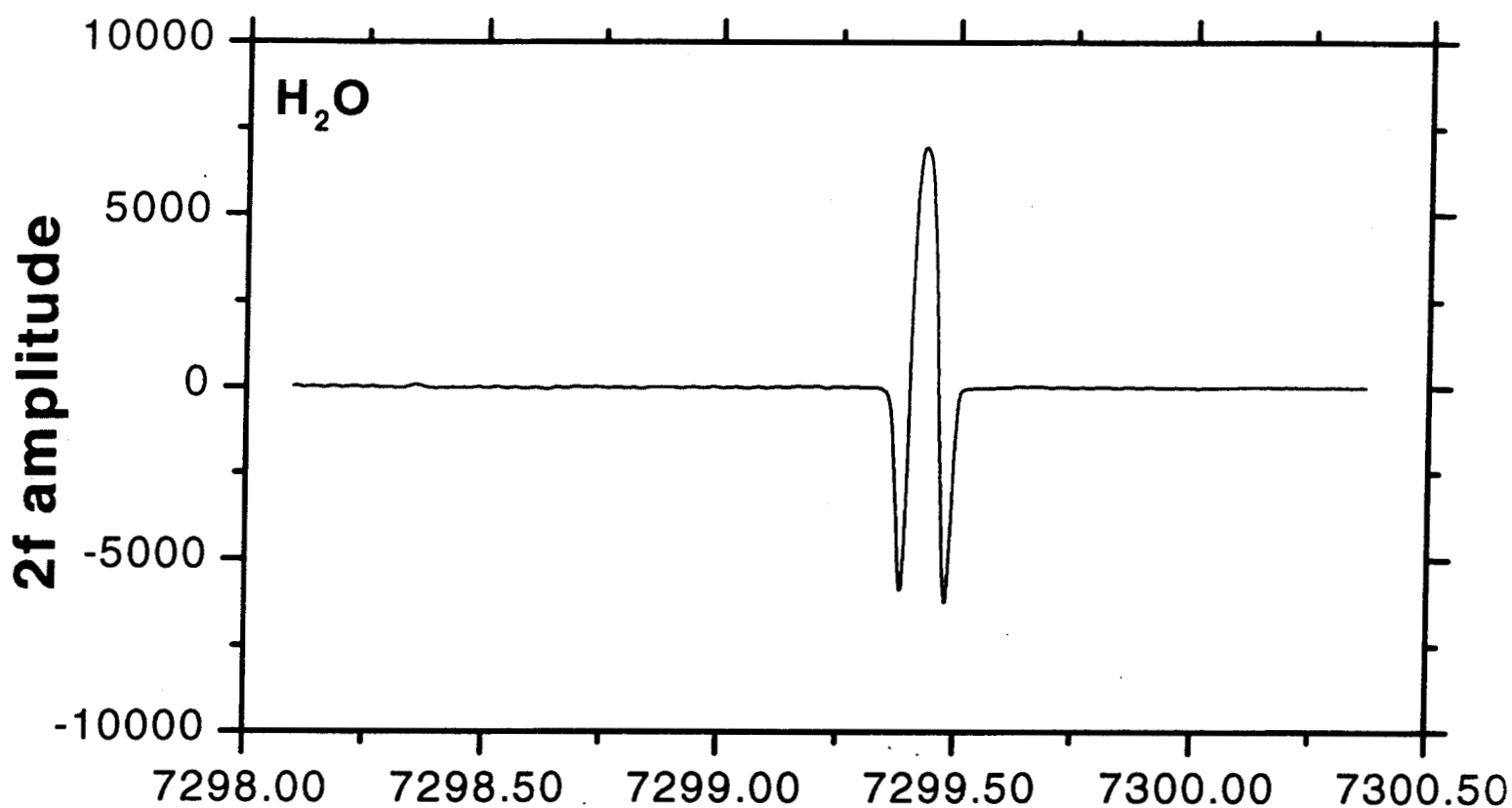
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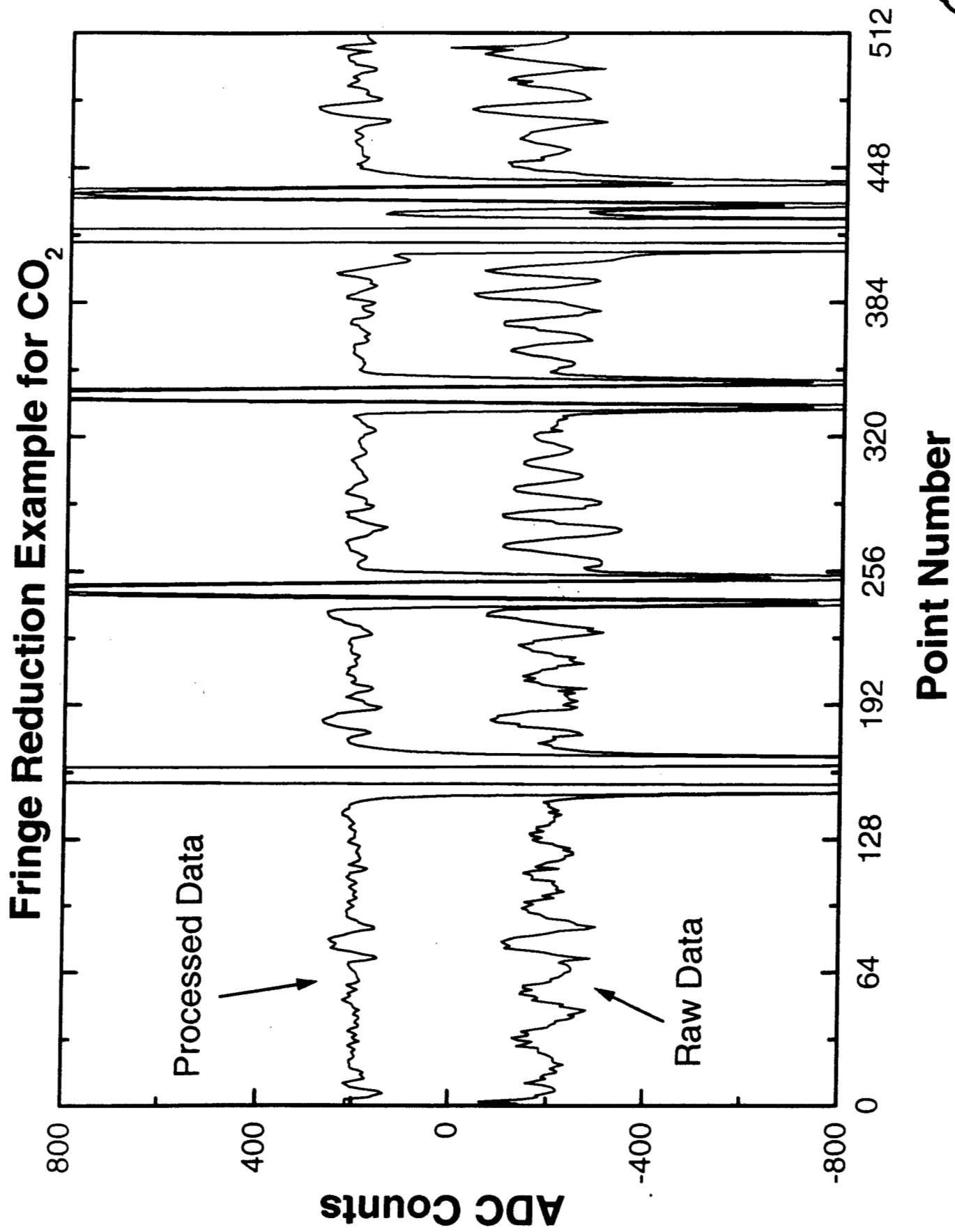
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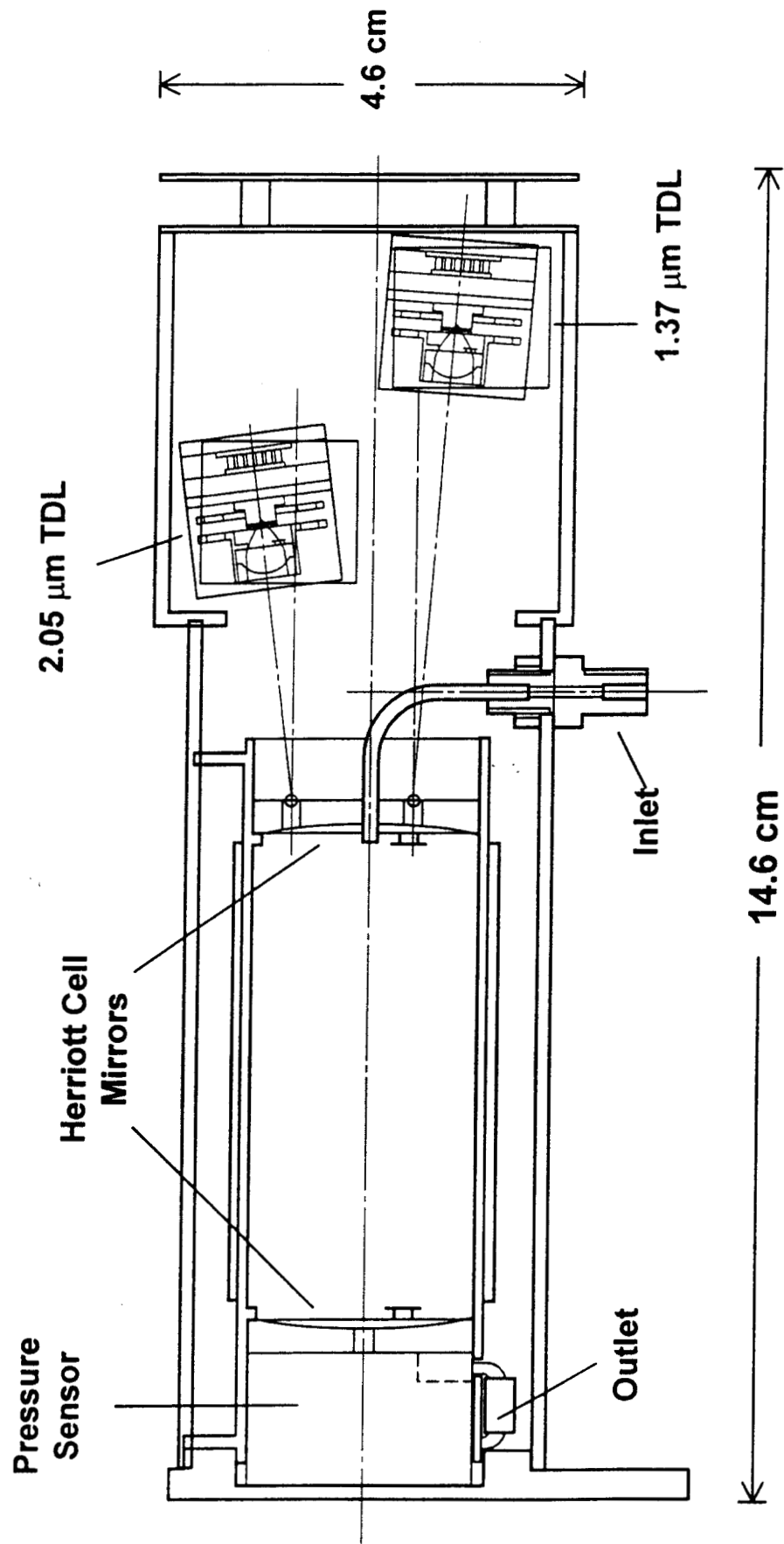


MET Spectral Regions





TEGA TDL Sensor (Plan View)



7

TEGA Spectral Regions

